Noise Correlations



Proposal:

E. Altman, E. Demler & M. Lukin PRA (2004) A. Polkovnikov et al., PNAS (2006) **Experiment:** Fölling et al., Nature (2005), Greiner et al., PRL (2005) Rom et al., Nature (2006)

Guarrera et al., PRL (2008)

related work:

Bach & Rzazewski, PRA (2004) Z. Hadzibabic et al. PRL (2004),

Yasuda & Shimizu, PRL (1996), Schellekens et al., Science (2005), Jeltes et al.,Nature (2007) Öttl et al., PRL (2005), Estève et al., PRL (2006),



Detecting Expanding Atom Clouds

Typically Noise in Images of a Mott Insulator

Single Image





200

0

400

Corrleations in Noise?



Hanbury-Brown Twiss effect correlates fluctuations at special distances r!

Quantitatively

 $g^{(2)}(r) - l > 0$ $g^{(2)}(r) - l = 0$ $g^{(2)}(r) - l < 0$ Noise correlated (Bosons)

Noise uncorrelated

Noise anti-correlated (Fermions)

- Hanbury Brown-Twiss Effect for Atoms (1) -



- Hanbury Brown-Twiss Effect for Atoms (2) -

There's another ways....



Hanbury Brown 1916-2002



- Hanbury Brown-Twiss Effect for Atoms (3) -

Cannot fundamentally distinguish between both paths...



Hanbury Brown 1916-2002



- Hanbury Brown-Twiss Effect for Atoms (4) -

Interference in Two-Particle Detection Probability!



- Multiple Wave Hanbury Brown-Twiss Effect (4) -

Interference in Two-Particle Detection Probability!



Deriving the Noise Correlation Signal (1)

In Time of Flight we measure:
$$\langle \hat{n}_{3D}(\mathbf{x}) \rangle_{\text{tof}} = \langle \hat{a}_{tof}^{\dagger}(\mathbf{x}) \hat{a}_{\text{tof}}(\mathbf{x}) \rangle_{\text{tof}}$$

 $\approx \langle \hat{a}^{\dagger}(\mathbf{k}) \hat{a}(\mathbf{k}) \rangle_{\text{trap}} = \langle \hat{n}_{3D}(\mathbf{k}) \rangle_{\text{trap}}$

where
$$\left(\mathbf{k} = M\mathbf{x}/\hbar t\right)$$

In Noise Correlations we measure:

$$\langle \hat{n}_{3D}(\mathbf{x}) \hat{n}_{3D}(\mathbf{x}') \rangle_{\text{tof}} \approx \langle \hat{a}^{\dagger}(\mathbf{k}) \hat{a}(\mathbf{k}) \hat{a}^{\dagger}(\mathbf{k}') \hat{a}(\mathbf{k}') \rangle_{\text{trap}} = \langle \hat{a}^{\dagger}(\mathbf{k}) \hat{a}^{\dagger}(\mathbf{k}') \hat{a}(\mathbf{k}') \hat{a}(\mathbf{k}) \rangle_{\text{trap}} + \delta_{\mathbf{k}\mathbf{k}'} \langle \hat{a}^{\dagger}(\mathbf{k}) \hat{a}(\mathbf{k}) \rangle_{\text{trap}} .$$

Deriving the Noise Correlation Signal (2)

$$\hat{a}(\mathbf{k}) = \int e^{-i\mathbf{k}\mathbf{r}} \hat{\psi}(\mathbf{r}) d^3 r$$
 with $\hat{\psi}(\mathbf{r}) = \sum_{\mathbf{R}} \hat{a}_{\mathbf{R}} w(\mathbf{r} - \mathbf{R})$

$$\implies \hat{a}(\mathbf{k}) = \tilde{w}$$

 $\tilde{v}(\mathbf{k})\sum_{\mathbf{R}}e^{-i\mathbf{k}\mathbf{R}}\hat{a}_{\mathbf{R}}$ Plug this into four operator correlator

For Mott state or Fermi gas, one has

$$\langle \hat{a}_{\mathbf{R}}^{\dagger} \hat{a}_{\mathbf{R}'} \rangle = n_{\mathbf{R}} \, \delta_{\mathbf{R},\mathbf{R}'}$$

which yields:

$$\langle \hat{n}_{3D}(\mathbf{x})\hat{n}_{3D}(\mathbf{x}')\rangle = |\tilde{w}(M\mathbf{x}/\hbar t)|^2 |\tilde{w}(M\mathbf{x}'/\hbar t)|^2 N^2$$
$$\times \left[1 \pm \frac{1}{N^2} \left|\sum_{\mathbf{R}} e^{i(\mathbf{x}-\mathbf{x}')\cdot\mathbf{R}(M/\hbar t)} n_{\mathbf{R}}\right|^2\right]$$

Information in the Noise – Correlations become visible!

$$g_{\exp}^{(2)}(\mathbf{b}) = \frac{\int \langle n(\mathbf{x} + \mathbf{b}/2) \cdot n(\mathbf{x} - \mathbf{b}/2) \rangle d^2 \mathbf{x}}{\int \langle n(\mathbf{x} + \mathbf{b}/2) \rangle \langle n(\mathbf{x} - \mathbf{b}/2) \rangle d^2 \mathbf{x}}$$



Fölling et al. Nature, 434, p. 481 (2005)



Let's change the sign...



Sympathetic Cooling of ⁴⁰K-⁸⁷Rb in Crossed Dipole Trap:



After final cooling in optical dipole trap 2×10⁵ ⁸⁷Rb (almost pure condensate) 2.5×10⁵ ⁴⁰K

After removal of ⁸⁷Rb

 $2 \times 10^{5} {}^{40}K @ T/T_F = 0.2$

Then load into 3D optical lattice and create a fermionic band insulator!

Adiabatic mapping: theory: A. Kastberg et al. PRL (1995) exp: M. Greiner et al., PRL (2001), M. Köhl et al. PRL (2005)



Mott insulator – Fermionic Band Insulator



Releasing the Fermi Gas



Noise Correlations of a Degenerate Fermi Gas



Rom et al. Nature **444**, 733 (2006)

First observation of fermionic antibunching for neutral atoms (maybe neutral particles)! (see also Jeltes et al., Nature 445, 402 (2007))

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An Alternative Description



Why Bosons and Fermions are Different in their Correlations



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Why Bosons and Fermions are Different in their Correlations



Now detection of many strongly correlated quantum states becomes possible!

Antiferromagnet

 $\langle \uparrow \rangle$

Spin wave

Charge density wave

1

In Situ Density Measurements

Absorption Imaging



Absorption per atom maximally approx 5-10% (typically smaller), mode-matching!

N. Gemelke et al., Nature 460, 995 (2009)

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State of the Art

D. Weiss, Penn State



Nelson et al., Nature Phys. 3

- fluorescence ir
- 5 μ m lattice sp



Detecting Atoms with an Electron Beam





P. Würtz et al., PRL (2009)

Optical Superlattices

See also Related Experiments at NIST (T. Porto & W. D. Phillips)

TTT

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Superimpose two standing waves with controllable phase & amplitude.



Array of double wells



17

Controllable Parameters

_______ • Intra & Interwell Barrier Depth

All parameters can be changed dynamically & in-situ!



Superlattices

Patterned loading of the short lattice



Patterned correlations



Detection of density wave via noise correlations.



2D Superlattice Geometries (1 SL)





LZ



Coupled Plaquette Systems

see B. Paredes & I. Bloch, PRA **77**, 23603 (2008) S. Trebst et al., PRL **96**, 250402 (2006)



Higher Lattice Orbital Physics

see V. Liu, A. Ho, C. Wu and others work exp: related to A. Hemmerich's exp.

