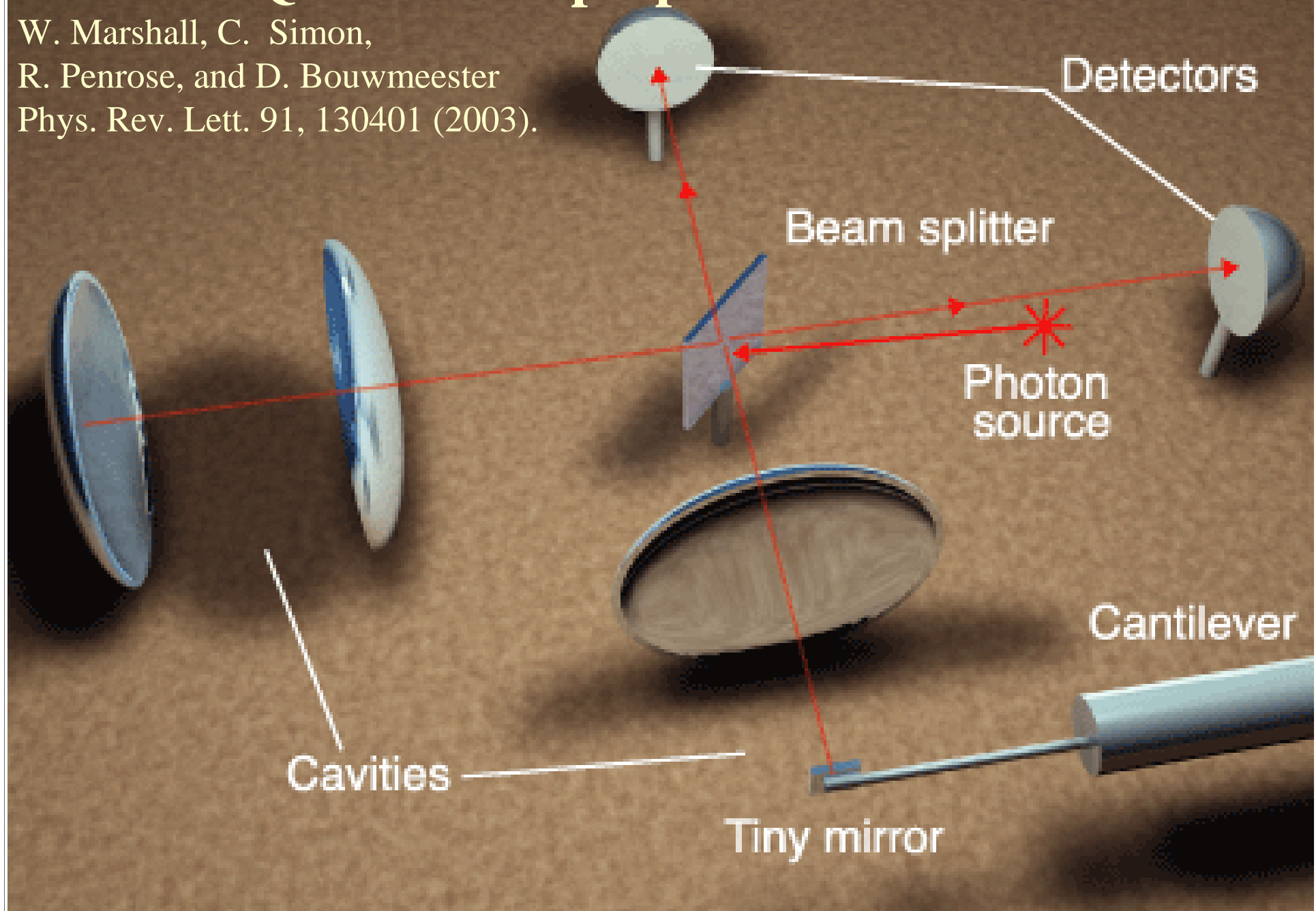
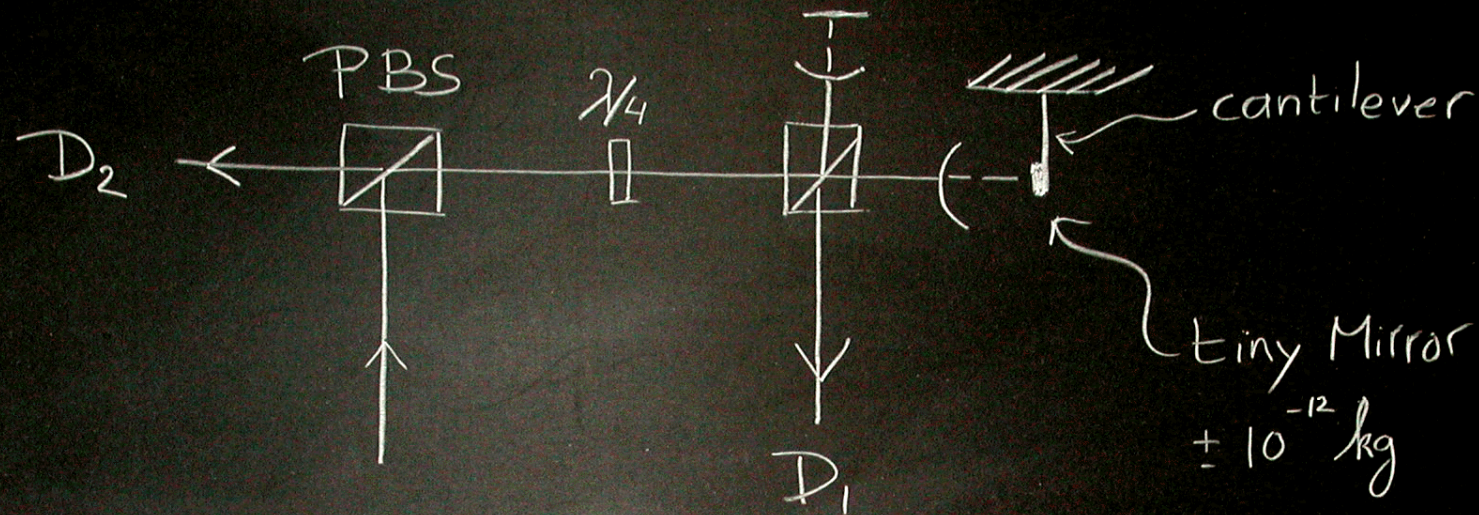


Towards Quantum Superpositions of a Mirror

W. Marshall, C. Simon,
R. Penrose, and D. Bouwmeester
Phys. Rev. Lett. 91, 130401 (2003).





$$\mathcal{H} = \hbar\omega_c a^\dagger a + \hbar\omega_m b^\dagger b - \hbar g a^\dagger a (b + b^\dagger)$$

$$g = \frac{\omega_c}{L} \sqrt{\frac{\hbar}{2M\omega_m}}$$

Law, PRA, **49**, 433 (1993)

Bose et al. PRA **59**, 3204 (1999)

Marshall et al. PRL **91**, 130401 (2003)

Optomechanical coupling $\kappa = (g/\omega_c)$

Consider single cavity mode: $\hat{H} = \hbar\omega_a \hat{a}^\dagger \hat{a}$ $\omega_a = \frac{n}{2}\omega_0$ $\omega_0 = 2\pi c/L$

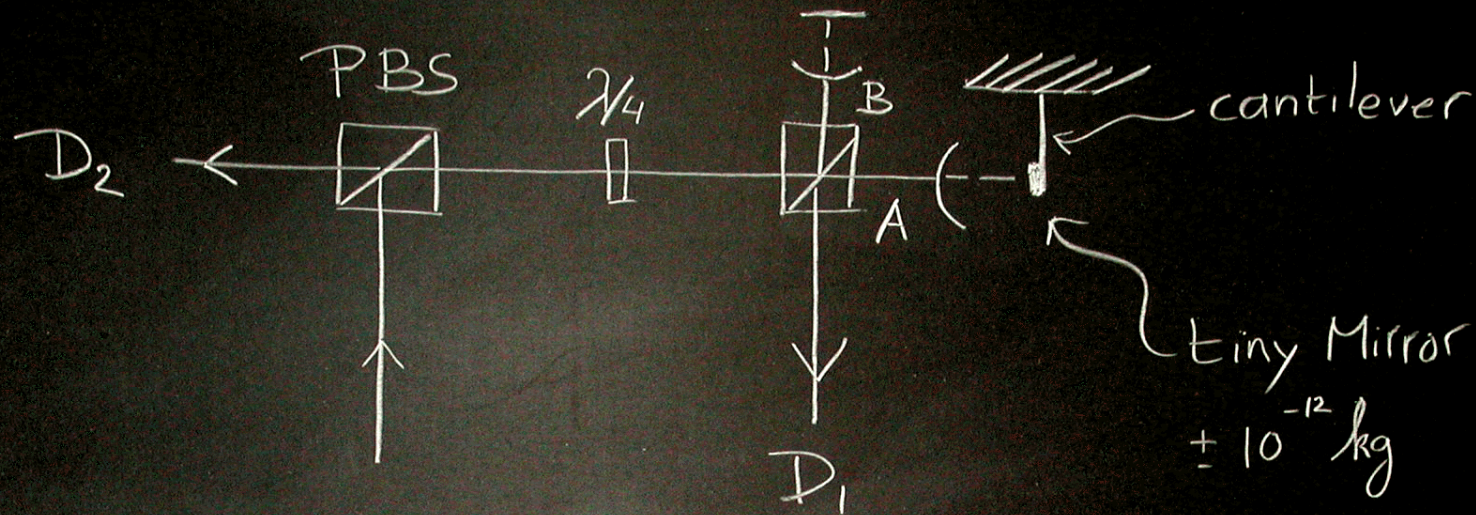
If mirrors moves by x : $\hat{H} = \hbar \frac{n}{2} \frac{2\pi c}{L+x} \hat{a}^\dagger \hat{a}$ $\hat{H} = \hbar \frac{n}{2} \omega_0 \frac{L}{L+x} \hat{a}^\dagger \hat{a}$
 $\approx \hbar\omega_a (1 - \frac{x}{L}) \hat{a}^\dagger \hat{a}$
 $= \hbar\omega_a \hat{a}^\dagger \hat{a} - \hbar \frac{\omega_a}{L} \hat{a}^\dagger \hat{a} x$

Quantize mirror position: $x \rightarrow \hat{x} = \sqrt{\hbar/2m\omega_c} (\hat{c} + \hat{c}^\dagger)$

Gives: $\hat{H} = \hbar\omega_a \hat{a}^\dagger \hat{a} - \hbar\omega_c \frac{\omega_a}{\omega_c L} \sqrt{\frac{\hbar}{2m\omega_c}} (\hat{c} + \hat{c}^\dagger) \hat{a}^\dagger \hat{a}$

$$\hat{H} = \hbar\omega_a \hat{a}^\dagger \hat{a} - \hbar\omega_c \kappa \hat{a}^\dagger \hat{a} (\hat{c} + \hat{c}^\dagger) \quad \kappa = \frac{\omega_a}{\omega_c} \frac{1}{L} \sqrt{\frac{\hbar}{2m\omega_c}}$$

Sorry on this his page ω_c is the cantilever frequency



Mirror in coherent state $|\beta\rangle = e^{-\frac{|\beta|^2}{2}} \sum_{n=0}^{\infty} \frac{\beta^n}{\sqrt{n!}} |n\rangle$

Initial state $|\psi(0)\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle_A |1\rangle_B + |1\rangle_A |0\rangle_B \right) |\beta\rangle$

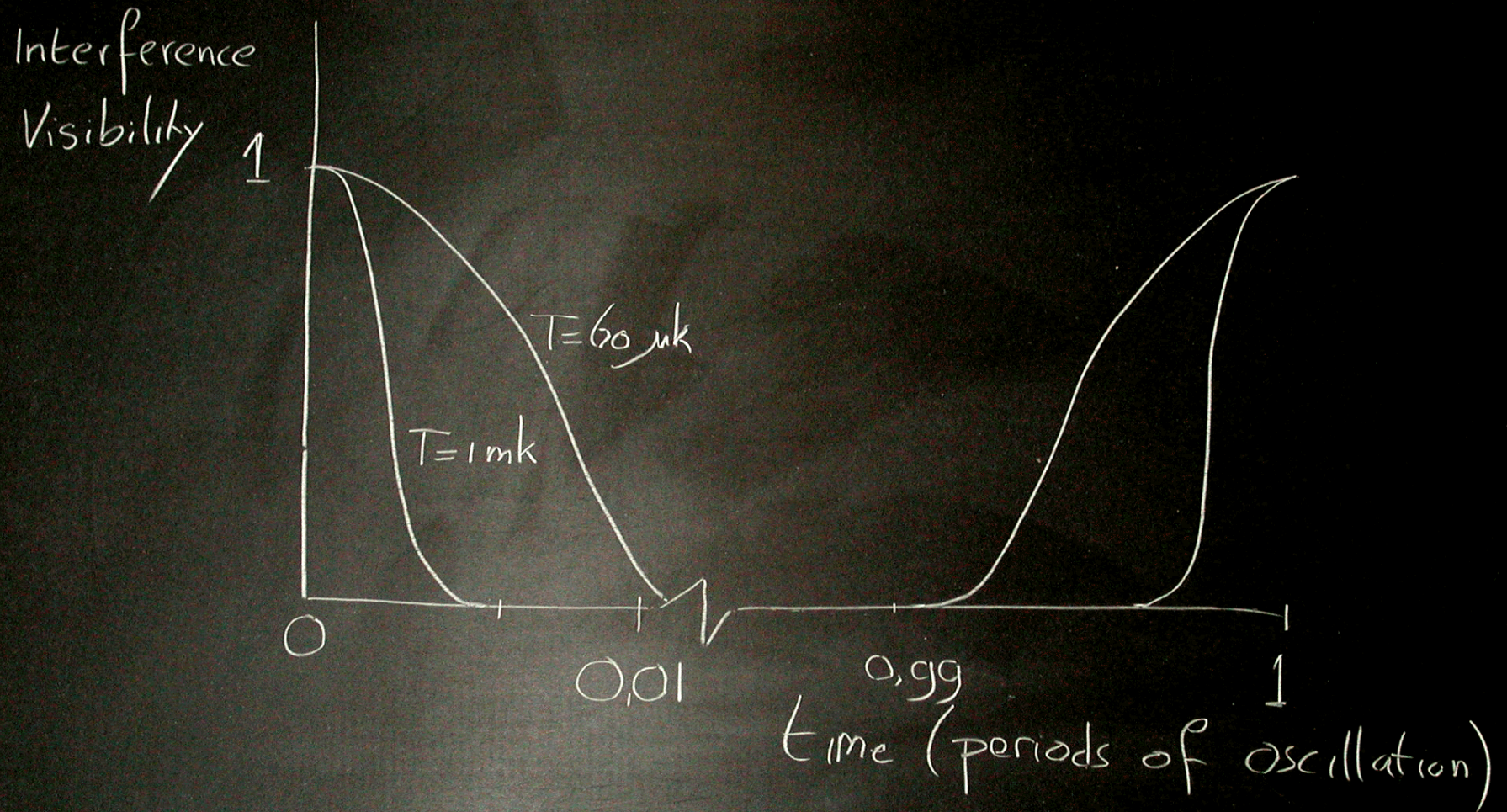
$|\psi(t)\rangle =$ entangled state of mirror and photon
except after full period of oscillation

$$\hat{U}(t) = e^{-i\hat{H}t/\hbar}$$

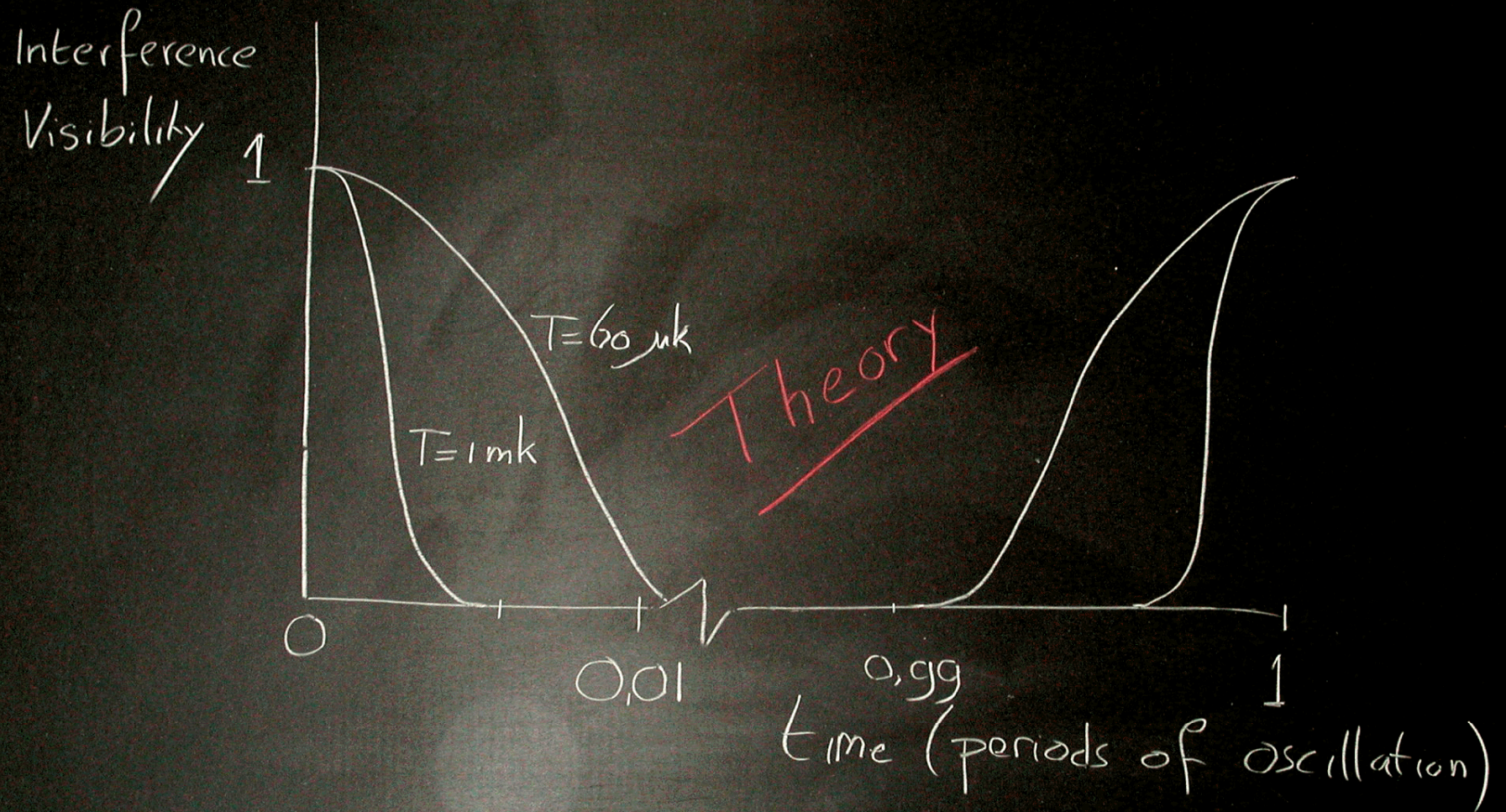
$$\hat{U}(t) = e^{-i\omega_a \hat{a}^\dagger \hat{a} t} e^{i\kappa^2 (\hat{a}^\dagger \hat{a})^2 (\omega_c t - \sin \omega_c t)} e^{\kappa \hat{a}^\dagger \hat{a} [\hat{c}^\dagger (1 - e^{-i\omega_c t}) - \hat{c} (1 - e^{i\omega_c t})]} e^{-i\omega_c \hat{c}^\dagger \hat{c} t}$$

$$\hat{U}|n\rangle|\beta\rangle = e^{-in\omega_a t} e^{i\kappa^2 n^2 (\omega_c t - \sin \omega_c t)} e^{i\kappa n \text{Im}[\beta(1 - e^{-i\omega_c t})]} |n\rangle |\beta e^{-i\omega_c t} + \kappa n(1 - e^{-i\omega_c t})\rangle$$

Note that after a full mechanical oscillation the cantilever is back in the same state as it was originally



T is effective temperature of the fundamental resonance of cantilever



T is effective temperature of fundamental resonance of cantilever

Overlap of two components of cantilever wave-function

$$|\langle \Phi_0(t) | \Phi_1(t) \rangle| = e^{-\kappa^2(1-\cos\omega_c t)}$$

0 after full oscillation and maximal after half periods

$$t = \pi/\omega_c \quad |\langle \Phi_0(\pi/\omega_c) | \Phi_1(\pi/\omega_c) \rangle| = e^{-2\kappa^2}$$

Requirement for creating superposition: $\kappa \gtrsim 1/\sqrt{2}$