

What is information

Represent content in binaries. 0 1

Communication



How can we make transmission most efficient?

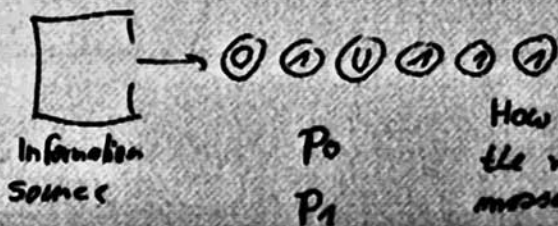
Extreme example I:

Digital camera picture of white sheet of paper. \rightarrow Represent content of the picture in a very compact form.

"All pixels are white"

Shannon 1948 or so, \rightarrow Information theory

For now: Each entry of the message is independent of the next.



P_0
 P_1

How much can I shorten the representation of the message without error

$$P_0 = \frac{1}{8}$$

$$P_1 = \frac{7}{8}$$

Enumerak

01111111	000
10111111	001
11011111	010
11101111	011
11110111	100
11111011	101
11111101	110
11111110	111



Length of sequences $\rightarrow \infty$

$P(\text{Typical sequence}) \rightarrow 1$

} Law of large numbers.

Given N digits: How many typical seq.?

$$\binom{N}{Np_0} = \frac{N!}{(Np_0)! (N(1-p_0))!} = \frac{N! e^{-N} (2\pi N)^{1/2}}{(Np_0)^{Np_0} e^{-Np_0} (2\pi Np_0)^{1/2} (N(1-p_0))^{N(1-p_0)} e^{-N(1-p_0)} (2\pi N(1-p_0))^{1/2}}$$

$k! \approx k^k e^{-k} \sqrt{2\pi k}$
Stirling for $n \rightarrow \infty$

$$= \frac{1}{\sqrt{2\pi N p_0}} \cdot \frac{N!}{N^{Np_0} p_0^{Np_0} N^{N(1-p_0)} (1-p_0)^{N(1-p_0)}} = \frac{1}{\sqrt{2\pi N p_0}} \cdot \left(\frac{1}{p_0^{p_0} (1-p_0)^{1-p_0}} \right)^N$$

$$\begin{aligned} \log_2 \binom{N}{N p_0} &\approx -\frac{1}{2} \log_2 (2^{2N} p_0^{p_0} p_1^{p_1}) + N \log_2 p_0^{p_0} p_1^{p_1} \\ &= -11 - \dots + N (p_0 \log p_0 + p_1 \log p_1) \\ &\approx N H(\{p_0\}) \end{aligned}$$

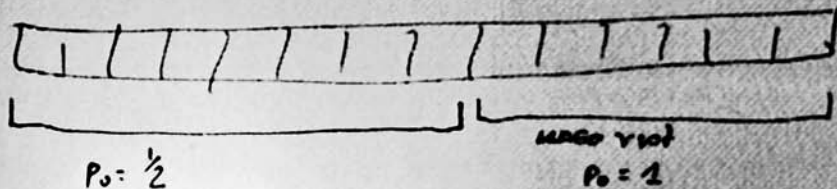
$$H(\{p_0\}) = -p_0 \log p_0 - p_1 \log p_1$$

↳ Shannon entropy

Can compress the original message by a factor $H(p_0)$.

Check: $p_0 = 0$ $H(p_0, p_1) = -0 \log 0 - 1 \log 1$
 $p_1 = 1$ $= 0 - 0$
 $= 0$

$p_0 = p_1 = \frac{1}{2}$ $H(p_0) = -\frac{1}{2} \log \frac{1}{2} - \frac{1}{2} \log \frac{1}{2}$
 $= -\log \frac{1}{2}$
 $= \log 2$
 $= 1$



$$P_0 P_{00} = 1$$

$$P_{01} = 0 = P_{10} = P_{11}$$

$0 \ 0$ uncorrelated
 $0 \ 0$ $I = 0$
 $0 \ 0$
 $0 \ 0$
 \vdots

Cryptography

Secret info

Secret correlation

0	⊕	0 = 0	⊕ 0 = 0
1	⊕	1 = 0	⊕ 1 = 1
1	⊖	1 = 0	⊖ 1 = 1
0	⊖	1 = 0	⊖ 1 = 1
1	⊕	0 = 0	⊕ 0 = 0
0	⊖	0 = 1	⊖ 0 = 1
1	⊕	1 = 1	⊕ 1 = 0

public
 communication

Secret correlations are a resource for
 secret public communication.

"One time pad"

Exercise: Transmit digit a via first entry of you "one time pad".

- 11 - b

- 11 -

Can any ~~one~~ learn something about a & b ?

Correlations

~~$H(X) = H(Y)$~~

Ex:

$$P_{00} = \frac{1}{2}$$

$$P_{01} = P_{10} = 0$$

$$P_{11} = \frac{1}{2}$$

$$\begin{matrix} 0 & 0 \\ 1 & 1 \\ 0 & 0 \\ 1 & 1 \\ 0 & 0 \\ 1 & 1 \\ 1 & 1 \\ \vdots & \vdots \\ 1 & 1 \end{matrix}$$

$$1 \ 1$$

$$0 \ 0$$

$$1 \ 1$$

$$0 \ 0$$

$$1 \ 1$$

$$1 \ 1$$

$$\vdots$$

$$1$$

By looking at first entry, I learn everything about the second entry.
 \Rightarrow maximally correlated.

~~$H(X, Y)$~~

entropy of joint

Random variable X, Y

Joint prob. distribution $p(x, y)$

$$I = H(X) - H(X|Y) = H(X) + H(Y) - H(X, Y)$$

exercice

$$H(X|Y) = \sum_y p(y) \sum_x -p(x|y) \log_2 p(x|y)$$

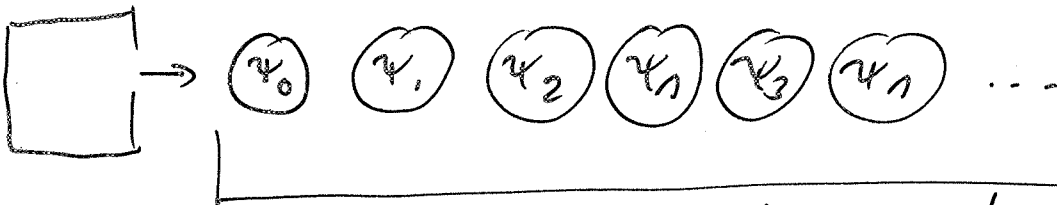
$$H(p(x)) = -\frac{1}{2} \log_2 \frac{1}{2} - \frac{1}{2} \log_2 \frac{1}{2} = 1$$

$$H(p(y)) = 1$$

$$H(p(x, y)) = 1$$

$$\Rightarrow I = 1 + 1 - 1 = 1$$

Quantum Info & Corr



N two level quantum systems

\Rightarrow Hilbert space is 2^N dimensional

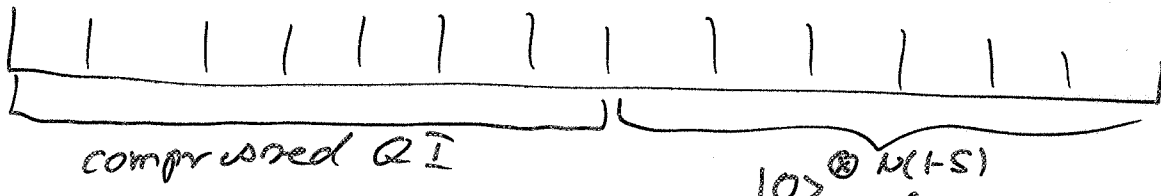
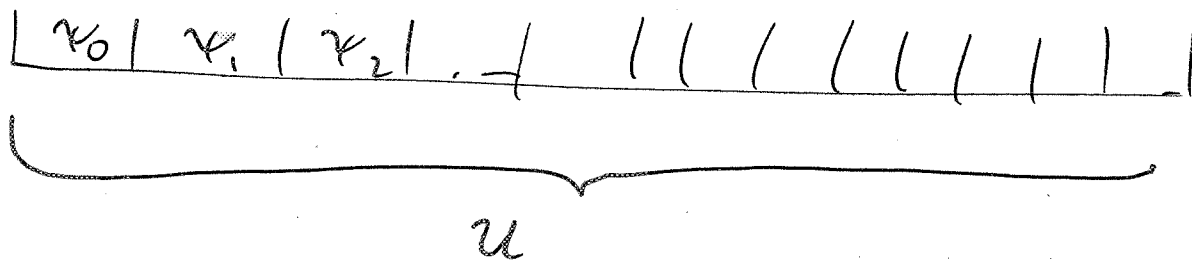
Quantum compression \rightarrow try and represent these quantum state in a smaller Hilbert space.

$$2^{NS}$$

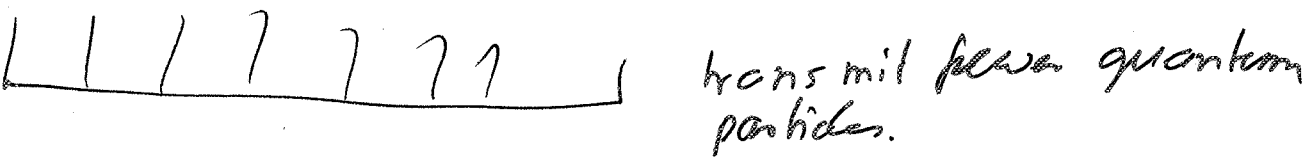
$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

$$= \sum_j p_j |\varphi_j\rangle\langle\varphi_j| \quad \langle\varphi_i|\varphi_j\rangle = \delta_{ij}$$

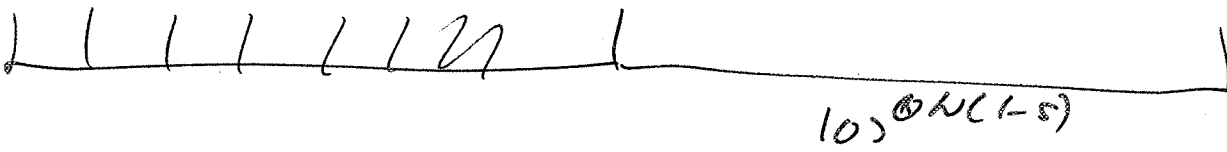
\rightarrow compression ratio $H(\rho) = \underline{\underline{S(\rho)}} = -\text{tr} \rho \log_2 \rho$



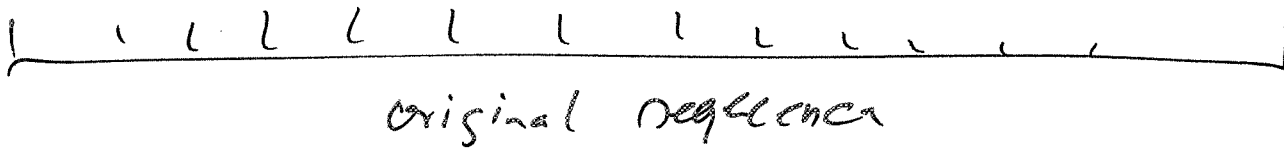
↓ throw away ↗



↓



↓ u^{-1}



typical sequences of quantum states

	n	
$ \varphi_0\rangle \varphi_1\rangle \varphi_2\rangle \varphi_1$	\rightarrow	$ \varphi_0\rangle \varphi_0\rangle 0\rangle 0\rangle$
$ \varphi_1\rangle \varphi_0\rangle \varphi_1\rangle \varphi_1$	\rightarrow	$ \varphi_0\rangle \varphi_1\rangle 0\rangle 0\rangle$
$ \varphi_1\rangle \varphi_1\rangle \varphi_0\rangle \varphi_1$	\rightarrow	$ \varphi_1\rangle \varphi_0\rangle 0\rangle 0\rangle$
$ \varphi_1\rangle \varphi_1\rangle \varphi_1\rangle \varphi_0$	\rightarrow	$ \varphi_1\rangle \varphi_1\rangle 0\rangle 0\rangle$

Schumacher 1995

Qubit

$$\rho = \frac{1}{2}I \Rightarrow S(\rho) = 1$$

Classical Info Shannon 1948

 atypical seq
typical seq

0001		0000
0010		0100
0100	→ enumerate	1000
1000		11 00

$$P_0 = 3/4$$

$$P_1 = 1/4$$

↓ compression by
a factor $H = -\sum_{i=0}^1 p_i \log_2 p_i$

total information content of a seq. of N digits
is $H \cdot N$
↳ Shannon entropy

Correlations

X Y
↗

measure X how much information do you
gain about Y

$$\begin{aligned} I &= H(X) - H(X|Y) \\ &= H(Y) - H(Y|X) \\ &= H(Y) + H(X) - H(X, Y) \end{aligned}$$

Quantum case

Schumacher 1995/95

$|\varphi_0\rangle |\varphi_1\rangle |\varphi_0\rangle \dots$

2-level systems Hilbert space of N systems has 2^N dimensions.

compresses quantum into smaller Hilbert space with dim $2^{N S(\rho)}$

where $\rho = \sum_i p_i |\varphi_i\rangle \langle \varphi_i|$

$$S(\rho) = -\text{tr} \rho \log_2 \rho \quad \text{von Neumann entropy}$$

Quantum State Teleportation

Alice is given an unknown quantum state of a qubit

$$\rho = \frac{1}{2} \cdot \underline{1}$$

If we can only do quantum operations locally in Ulm (Alice's lab) and in Strathelyde (Bob's lab) & they can telephone to correlate their actions.



$$\begin{aligned}
 (\alpha|00\rangle + \beta|11\rangle)^{\otimes 2} &= (\alpha|0_A\rangle|0_B\rangle + \beta|1_A\rangle|1_B\rangle) \otimes (\alpha|0_A\rangle|0_B\rangle + \beta|1_A\rangle|1_B\rangle) \\
 &= \alpha^2|00\rangle_A|00\rangle_B \\
 &\quad + \alpha\beta|01\rangle_A|01\rangle_B \\
 &\quad + \alpha\beta|10\rangle_A|10\rangle_B \\
 &\quad + \beta^2|11\rangle_A|11\rangle_B
 \end{aligned}$$

Alice measures how many of her particles are in state $|1\rangle$.

$$\begin{aligned}
 A &= 0 \cdot |00\rangle \times |00\rangle + 1 \cdot (|01\rangle \times |01\rangle + |10\rangle \times |10\rangle) \\
 &\quad + 2|11\rangle \times |11\rangle
 \end{aligned}$$

If Alice find eigenvalue 1 \Rightarrow

$$u \otimes u \quad \frac{|01\rangle|01\rangle + |10\rangle|10\rangle}{\sqrt{2}} \quad \text{in state after measurement}$$

Alice and Bob change their basis locally

$$\begin{aligned}
 |00\rangle &\rightarrow |00\rangle \\
 |01\rangle &\rightarrow |01\rangle \\
 |10\rangle &\rightarrow |11\rangle \\
 |11\rangle &\rightarrow \cancel{|11\rangle} |10\rangle
 \end{aligned}$$

$$u = |00\rangle \times |00\rangle + |01\rangle \times |01\rangle + |11\rangle \times |10\rangle + |10\rangle \times |11\rangle$$

$$uu^\dagger = u^\dagger u = \mathbb{1}$$

$$\frac{|01\rangle|01\rangle + |11\rangle|11\rangle}{\sqrt{2}} = \frac{|0\rangle_A|1\rangle_B + |1\rangle_A|1\rangle_B}{\sqrt{2}} \quad |1\rangle_A|1\rangle_B$$

$$(\alpha|00\rangle + \beta|11\rangle)^{\otimes N}$$

\Rightarrow transform this with unit probability ($N \rightarrow \infty$) to $N \cdot H(\alpha^2)$ perfectly entangled pairs.

\Rightarrow The state $\alpha|00\rangle + \beta|11\rangle$ contains $H(\alpha, \beta)$ ebits.

\square

$$\rho_A = \alpha^2 |0\rangle\langle 0| + \beta^2 |1\rangle\langle 1|$$

$$H(\alpha, \beta) = S(\rho_A)$$

$$p_0 = \alpha^2$$

$$p_1 = \beta^2$$

In For mixed states I could define the amount of entanglement as that required to transfer 1 qubit of quantum information.

\Rightarrow How many $|\phi^+ \rangle \langle \phi^+|$ can I get from $\rho^{\otimes N}$?

$$|\phi^+ \rangle \langle \phi^+| \rightarrow \rho \quad \rho \rightarrow |\phi^+ \rangle \langle \phi^+|$$

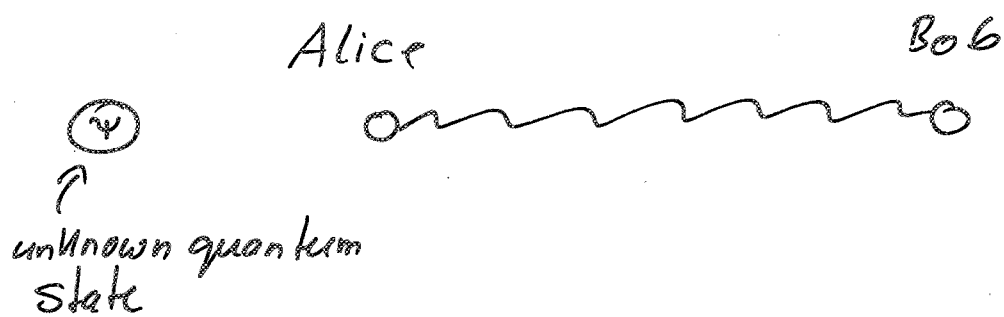
=

How many $|\phi^+ \rangle \langle \phi^+|$ do I need to make $\rho^{\otimes N}$ using local operations and classical correlations?

$$|\phi^+ \rangle \langle \phi^+| \rightarrow \rho$$

Both are valid ways to quantify correlations but they lead to different answers.

Alice & Bob need additional resource.



$$|\phi^\pm\rangle = \frac{|00\rangle \pm |11\rangle}{\sqrt{2}}$$

$$|\psi^\pm\rangle = \frac{|01\rangle \pm |10\rangle}{\sqrt{2}}$$

$$(\underbrace{\alpha|0\rangle_A + \beta|1\rangle_A}_{|\psi\rangle}) \otimes \frac{|0\rangle_A |0\rangle_B + |1\rangle_A |1\rangle_B}{\sqrt{2}}$$

$$= \alpha|00\rangle_A |0\rangle_B + \alpha|01\rangle_A |1\rangle_B + \beta|10\rangle_A |0\rangle_B + \beta|11\rangle_A |1\rangle_B$$

$$= \alpha(|\phi^+\rangle + |\phi^-\rangle)|0\rangle_B + \alpha(|\psi^+\rangle + |\psi^-\rangle)|1\rangle_B + \beta(|\psi^+\rangle - |\psi^-\rangle)|0\rangle_B + \beta(|\phi^+\rangle - |\phi^-\rangle)|1\rangle_B$$

$$= |\phi^+\rangle_A (\alpha|0\rangle_B + \beta|1\rangle_B) + |\phi^-\rangle_A (\alpha|0\rangle_B - \beta|1\rangle_B) + |\psi^+\rangle_A (\alpha|1\rangle_B + \beta|0\rangle_B) + |\psi^-\rangle_A (\alpha|1\rangle_B - \beta|0\rangle_B)$$

$$+ |\psi^+\rangle_A (\alpha|1\rangle_B + \beta|0\rangle_B) + |\psi^-\rangle_A (\alpha|1\rangle_B - \beta|0\rangle_B)$$

$$+ |\psi^+\rangle_A (\alpha|1\rangle_B + \beta|0\rangle_B) + |\psi^-\rangle_A (\alpha|1\rangle_B - \beta|0\rangle_B)$$

$$+ |\psi^-\rangle_A (\alpha|1\rangle_B - \beta|0\rangle_B)$$

$$= |\phi^+\rangle_A (\alpha|0\rangle_B + \beta|1\rangle_B) + |\phi^-\rangle_A \sigma_z (\alpha|0\rangle_B + \beta|1\rangle_B) + |\psi^+\rangle_A \sigma_x (\alpha|0\rangle_B + \beta|1\rangle_B) + |\psi^-\rangle_A \sigma_z \sigma_x (\alpha|0\rangle_B + \beta|1\rangle_B)$$

$$+ |\phi^-\rangle_A \sigma_z (\alpha|0\rangle_B + \beta|1\rangle_B)$$

$$+ |\psi^+\rangle_A \sigma_x (\alpha|0\rangle_B + \beta|1\rangle_B)$$

$$+ |\psi^-\rangle_A \sigma_z \sigma_x (\alpha|0\rangle_B + \beta|1\rangle_B)$$

$$\sigma_z = |0\rangle\langle 0| - |1\rangle\langle 1|$$

$$\sigma_x = |0\rangle\langle 1| + |1\rangle\langle 0|$$

$$A = 1|\phi^+ \times \phi^+| + 2|\phi^- \times \phi^-| + 3|\psi^+ \times \psi^+| + 4|\psi^- \times \psi^-|$$

Alice makes measurements of A and
she communicates the outcome to Bob.

Ex: Alice finds "2" \rightarrow projects onto $|\phi^-\rangle$

\rightarrow Bob's particles is in state

$$|\phi^-\rangle_A \otimes \sigma_z (\alpha |0\rangle_B + \beta |1\rangle_B)$$

Bob applies σ_z to his particle

$$|\phi^-\rangle \otimes \sigma_z (\sigma_z (\alpha |0\rangle + \beta |1\rangle))$$

$$= |\phi^-\rangle \otimes (\alpha |0\rangle + \beta |1\rangle)$$

$$\sigma_x = \sigma_x^\dagger \quad \sigma_x \sigma_x^\dagger = \sigma_x \sigma_x = \mathbb{1}$$

$$\sigma_x^\dagger \sigma_x = \mathbb{1}$$

Deterministic process

Bennett et al PRL 1993

Transmits quantum information

$$\square \otimes \otimes \otimes (\phi_0) \dots$$

$$\rho = \frac{1}{2} \mathbb{1} \quad \text{"unknown state"}$$

$$S(\rho) = 1$$

1 unit of quantum correlation^(1 ebit) are necessary and sufficient to teleport one qubit of quantum information.

$$\frac{|00\rangle + |11\rangle}{\sqrt{2}} = 1 \text{ ebit}$$

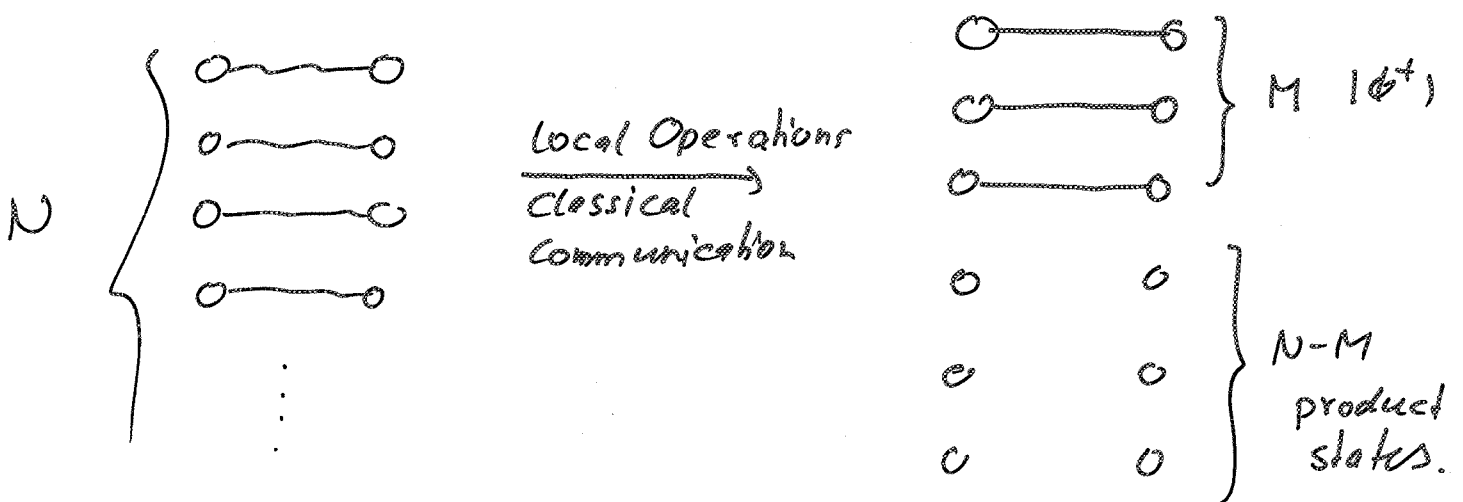
Question: How ~~many~~ much quantum correlations are contained in

$$\alpha |00\rangle + \beta |11\rangle ?$$

$$\alpha = \beta = \frac{1}{\sqrt{2}} \Rightarrow 1 \text{ ebit}$$

$$\begin{matrix} \alpha = 0 \\ \beta = 1 \end{matrix} \Rightarrow |1\rangle|1\rangle \Rightarrow 0 \text{ ebit}$$

$$(\alpha |00\rangle_{AB} + \beta |11\rangle_{AB})^{\otimes N}$$



If true for all $\rho \Delta \sigma$

$$E_1(\rho) > E_1(\sigma)$$

$$E_2(\rho) > E_2(\sigma)$$

$$\Rightarrow \bar{E}_1 \equiv E_2$$

Classical conditions

$$I = H(X) + H(Y) - H(X, Y)$$

□

$$I = S(\rho_A) + S(\rho_B) - S(\rho_{AB})$$

$$|\phi^+ \otimes \phi^+\rangle = \rho_{AB}$$

$$S(\rho_{AB}) = 0$$

$$S(\rho_A) = 1$$

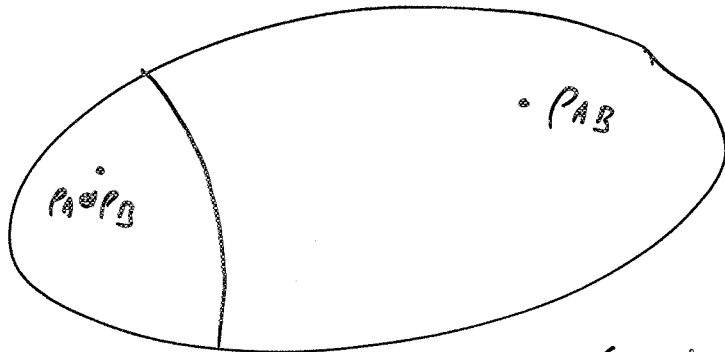
$$S(\rho_B) = 1$$

$$I = 2$$

$$I = S(\rho_{AB} \parallel \rho_A \otimes \rho_B)$$

$$S(\sigma \parallel \rho) = \text{tr}(\sigma \log \sigma - \sigma \log \rho)$$

relative entropy



$$E(\rho_{AB}) = \min_{\sigma \in \Sigma(\rho; \sigma_A, \sigma_B)} S(\rho_{AB} \parallel \sigma)$$

$$\sum_{i \in \Lambda} p_i \rho_i^A \otimes \rho_i^B$$

Plenio & Virmani, Quant. Inf. Comp. 7, 1 (2007)
~~arXiv:05~~

~~Vedral~~
Plenio & Vedral, Cont. Phys. 39, 431 (1998)

Horodecki^{Ⓢ4}, Rev Mod Phys. ... (2009)